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Appendix

Method for solving the Fredholm integral equation of first kind numerically

Consider a Fredholm integral equation of first the kind as (Delves and Mohamed (1985))

$$\int_a^b N(z, x, s)\psi(z, s)dz = L(x, s) \quad (\text{A.1})$$

By following Delves and Mohamed (1985) and Sherief and El-Maghreby (2005), Eq. (A.1) can be rewritten as

$$N\psi = L \quad (\text{A.2})$$

where $L(x, s)$ will be known to some accuracy ϵ . Therefore, the solution of Eq. (A.2) is expected to satisfy the accuracy constraint $\|N\psi - L\| \leq \epsilon$. Now, the objective is to find a smoothest function (ψ) satisfying Eq. (A.2) and the accuracy constraint such that for some linear operator F , the norm of $F\psi$ has minimum value. This objective leads to the following constrained minimization problem

$$\min_{\psi} \|F\psi\|, \quad \text{subjected to, } \|N\psi - L\| \leq \epsilon. \quad (\text{A.3})$$

Deriving solution to the minimization problem Eq. (A.3) is possible for any given norm, still, it is relatively complicated to solve the problem analytically. Therefore, instead of solving this problem, Delves and Mohamed (1985) have solved a related problem numerically.

From the present problem we can observe that for some operator, F the minimum of

$\|F\psi\|$ will decrease with increase in ϵ , i.e. as the constrained is weakened. Therefore, for the minimization of Eqs. (A.1-A.2), the constrain will be binding to $\|N\psi - L\| = \epsilon$. Now, by solving the unconstrained problem

$$\min_{\psi} \|N\psi - L\|^2 + \kappa \|F\psi\|^2, \quad (\text{A.4})$$

we will find a minimum value, ς of $\|N\psi - L\|$. Now, $\kappa \rightarrow 0$, $\varsigma \rightarrow 0$ implies that the solution of the Eqs. (A.1) and (A.2) exist and for some value of κ , $\varsigma = \epsilon$. Therefore, the solution of the Eq. (A.4) is identical to the solution of Eqs. (A.1-A.2) (see Delves and Mohamed (1985)). Since Eq. (A.4) is an unconstrained problem therefore, it is easy to solve this problem as compared to Eqs. (A.1-A.2). This unconstrained problem is referred to as regularization problem, and the numerical method used for solving this problem is known as regularization method (see refs. Delves and Mohamed (1985); Sherief and El-Magharby (2005)).

The most common choice for the operator F is $1, \frac{d}{dx}, \frac{d^2}{dx^2}$. Here, we take $F = 1$ along with the usual \mathbb{L}^2 norm for computation purpose.

Now, the Eq. (A.4) will takes the form

$$\min_{\psi} \mathcal{M}(\psi) = \langle N\psi - L, N\psi - L \rangle + \kappa \langle \psi, \psi \rangle \quad (\text{A.5})$$

Here (\cdot) denote the scalar product in \mathbb{L}^2 norm.

Rearranging Eq. (A.5), we find

$$\mathcal{M}(\psi) = \langle \psi, \{N^+N + \kappa I\}\psi \rangle - \langle \psi, N^+L \rangle - \langle N^+L, \psi \rangle + \langle L, L \rangle \quad (\text{A.6})$$

where N^+ denote the Hermitian conjugate of N .

Therefore, for the minimum value of \mathcal{M} at any point ψ for any arbitrary function g , we have

$$\left. \frac{\partial \mathcal{M}(\psi + \epsilon g)}{\partial \epsilon} \right|_{\epsilon=0} = 0 \quad (\text{A.7})$$

Now, in view of Eqs. (A.6) and (A.7), we find the relation for minimum ψ as

$$(N^+N + \kappa I)\psi = N^+L \quad (\text{A.8})$$

which implies that

$$\int_a^b \hat{N}(z, x, s)\psi(z, s)dz + \kappa\psi(z, s) = \hat{L}(x, s) \quad (\text{A.9})$$

where

$$\hat{N}(z, x, s) = \int_a^b N^*(z, u, s)N(u, x, s)du \quad (\text{A.10})$$

$$\hat{L}(x, s) = \int_a^b N^*(z, x, s)L(z, s)dz \quad (\text{A.11})$$

Here $N^*(z, x, s)$ is the complex conjugate of $N(z, x, s)$.

Thus, the Fredholm integral equation of first kind (A.1) is now transformed into a Fredholm integral equation of the second kind with iterated kernels given by Eq. (A.9), and the parameter κ is taken to be 10^{-5} (see Delves and Mohamed (1985)).

Next, we present a numerical method for solving a Fredholm integral equation of the second kind. For this purpose, we consider an integral equation of the form

$$\Psi(\varphi, s) + \int_a^b N(\varphi, u, s)\Psi(u, s)du = L(\varphi, s) \quad (\text{A.12})$$

Eq. (A.12) can be approximated as

$$\Psi(\varphi, s) + \sum_{j=0}^m S_j N(\varphi, u_j, s)\Psi(u_j, s) \approx L(\varphi, s) \quad (\text{A.13})$$

where $u_j, j = 0, 1, 2, \dots, m$ are the $m + 1$ equally spaced points in the interval $[a, b]$, and S_j denotes the corresponding weight constant. Now, we assume that the Eq. (A.13) satisfies the chosen points $u_j, j = 0, 1, 2, \dots, m$. Therefore, we find the system of $m + 1$ linear equations as

$$\Psi(u_i, s) + \sum_{j=0}^m S_j N(u_i, u_j, s)\Psi(u_j, s) \approx L(u_i, s), \quad i = 0, 1, 2, \dots, m, \quad (\text{A.14})$$

in the $m + 1$ unknowns $\Psi(u_0, s), \Psi(u_1, s), \dots, \Psi(u_m, s)$, that specify the approximation of the function $\Psi(u, s)$ at $m + 1$ chosen points.

Now, if we take the notation $\Psi(u_i, s) = \Psi_i$, $L(u_i, s) = L_i$, and $N(u_i, u_j, s) = N_{ij}$, $i, j = 0, 1, \dots, m$, then from equation (A.14) we have

$$\Psi_i + \sum_{j=0}^m S_j N_{ij}\Psi_j \approx L_i, \quad i = 0, 1, \dots, m, \quad (\text{A.15})$$

Therefore, the system of Eq. (A.15) can be written in the matrix form as

$$\boldsymbol{\Psi} + \mathbf{NS}\boldsymbol{\Psi} = \mathbf{L}$$

or,

$$(\mathbf{I} + \mathbf{NS})\boldsymbol{\Psi} = \mathbf{L} \tag{A.16}$$

where Ψ_i and L_i are the components of vector $\boldsymbol{\Psi}$ and \mathbf{L} , respectively and matrices \mathbf{N} and \mathbf{S} are defined as $\mathbf{N} = [N_{ij}]$, $\mathbf{S} = [S_i\delta_{ij}]$, \mathbf{I} is the unity matrix of order $m + 1$.

For the numerical computation, the weighting constants S_j are taken by following Simpson's one-third rule of integration as

$$\begin{aligned} S_0 = S_m &= \frac{h}{3}, \\ S_{2i-1} &= \frac{4h}{3}, \quad S_{2i} = \frac{2h}{3}, \quad i = 1, 2, 3, \dots, \frac{m}{2} \end{aligned} \tag{A.17}$$

PUBLICATIONS AND CONFERENCES

Publication Related to the Thesis:

- **Om Namha Shivay**, and Santwana Mukhopadhyay. “A porothermoelasticity theory for anisotropic medium.” *Continuum Mechanics and Thermodynamics* (2021): 1-18. DOI. 10.1007/s00161-021-01030-2
- **Om Namha Shivay**, and Santwana Mukhopadhyay. “Variational principle and reciprocity theorem on the temperature-rate-dependent poro-thermoelasticity theory.” *Acta Mechanica* (2021): 1-13. DOI. 10.1007/s00707-021-02996-5
- **Om Namha Shivay**, and Santwana Mukhopadhyay. “A complete Galerkin’s type approach of finite element for the solution of a problem on modified Green–Lindsay thermoelasticity for a functionally graded hollow disk.” *European Journal of Mechanics-A/Solids* 80 (2020): 103914.
- **Om Namha Shivay**, and Santwana, Mukhopadhyay. “On the temperature-rate dependent two-temperature thermoelasticity theory.” *Journal of Heat Transfer* 142.2 (2020): 022102.
- **Om Namha Shivay**, and Santwana Mukhopadhyay. “On the solution of a problem of extended thermoelasticity theory (ETE) by using a complete finite element approach.” *Computational Methods in Science and Technology* 2 (2019): 5.
- **Om Namha Shivay**, and Santwana Mukhopadhyay. “An investigation on modified temperature-rate dependent two-temperature thermoelasticity theory.” (Communicated to an International Journal)

Publications Apart from Thesis:

- **Om Namha Shivay**, and Santwana Mukhopadhyay. “Some basic theorems on a recent model of linear thermoelasticity for a homogeneous and isotropic medium.” **Mathematics and Mechanics of Solids** 24.8 (2019): 2444-2457.
- Anil Kumar, **Om Namha Shivay**, and Santwana Mukhopadhyay. “Infinite speed behavior of two-temperature Green–Lindsay thermoelasticity theory under temperature-dependent thermal conductivity.” **Zeitschrift für angewandte Mathematik und Physik** 70.1 (2019): 1-16.
- Shashi Kant, Manushi Gupta, **Om Namha Shivay**, and Santwana Mukhopadhyay. “An investigation on a two-dimensional problem of Mode-I crack in a thermoelastic medium.” *Zeitschrift für angewandte Mathematik und Physik* 69.2 (2018): 21.

Conferences and Workshops:

- Participated in 64th International Congress of ISTAM held at IIT Bhubaneswar during December 9-12, 2019 and presented a work with the title “An investigation of the Moore-Gibson-Thompson thermoelasticity theory using a meshfree finite element method.”
- Participated in International Conference on Differential Equations and Control Problems: Modeling, Analysis and Computations (ICDECP19) held at IIT Mandi during June 17-19, 2019 and presented a work with the title “Finite element solution of a problem on coupled thermoelasticity for functionally graded material using two different approaches for time domain.”
- Participated in QIP short term Course on Computational Methods on Integral and differential Equations held at IIT (BHU), Varanasi, during December 10-16, 2018.
- Participated in International Conference on Engineering, Computers and Natural Sciences 2018 held at Vivanta by Taj, Panjim, Goa during October 19-21, 2018 and presented the paper with the title “Application of complete finite element method on a coupled thermoelasticity theory of type III.”
- Participated in GIAN Course on Fractional Derivatives and Its Applications held at IIT (BHU), Varanasi, during January, 30- February 03, 2018.
- Attended training program on Tools for Scientific Documentation: LaTeX, JabRef, DocEar and other open source software held at DST, Banaras Hindu University during January 2-12, 2018.

- Attended a workshop on “Smart Materials and Structures – Recent Trends in Industrial Applications” held at IIT (BHU) during September 4-11, 2017.
