# Computing energy budget within a crop canopy from Penmann's formulae

 ${\rm MAHENDRA}\ {\rm MOHAN^*}$  and  ${\rm K}\ {\rm K}\ {\rm SRIVASTAVA^{**}}$ 

\*Radio and Atmospheric Science Division, National Physical Laboratory, New Delhi 110012, India \*\*Department of Chemical Engineering, Institute of Technology, Banaras Hindu University, Varanasi 221005, India

The Lhomme's model (1988a), that extended Penmann's formulae to a multi-layer model, is redefined as a function of micrometeorological and physiological profiles of crop canopy. The sources and sinks of sensible and latent heat fluxes are assumed to lie on a fictitious plane called zerodisplacement plane. Algorithms are given to compute sensible and latent heat flux densities. Performance of the algorithms is compared with that of earlier algorithms.

# 1. Introduction

For natural surfaces, the partitioning of available energy (Rn - S) into sensible heat flux (H) and latent heat flux  $(\lambda E)$  is accounted by the well-known Penmann's equations (Penmann 1948, 1953):

$$H = \frac{\Delta(Rn - S) - \rho c_p Da/r_a}{\Delta + \gamma (1 + r_s/r_a)},$$
 (1)

$$\lambda E = \frac{\Delta (Rn - S) + \rho c_p Da/r_a}{\Delta + \gamma (1 + r_s/r_a)},$$
 (2)

where Rn is the net radiation, S is the soil heat flux, Da is the vapour deficit of air at a reference height,  $\gamma$  is the psychrometric constant,  $\Delta$  is the slope of saturated vapour pressure curve at the air temperature,  $c_p$  the specific heat of air at constant pressure, and  $\rho$  the mean density of air.  $r_a$  is the aerodynamic resistance calculated between the surface level and the reference height and  $r_s$  being the surface resistance for water vapour transfer. In practice, there are two types of models: single layer and multi-layer models. The single layer model describes the convective transfers from a surface comprising of vegetation and underlying surface as a lumped system. Such a surface is treated as a single source. Monteith (1981) applied this approach to a vegetation stand on the assumption that the vegetation stand acts as a single equivalent surface absorbing radiative energy and transfering sensible and latent heat fluxes to the air. Deardorff (1978)

developed single layer parameterization of vegetation that involved the solution of energy budget equation.

In multi-layer models (Waggoner and Reifsnyder 1968; Shuttleworth 1976; Chen 1984; Meyers and Paw 1986; and Naot and Mahrer 1989), vertical transport of sensible and latent heat is described by considering a continuous or discrete set of horizontal planes, each one exchanging heat and vapour with the air. These models include parameterization of stomatal response to the environmental conditions, along with some models for heat and water fluxes into and out of soil layer.

Lhomme (1988a) presented a model based on electrical analogue that simulates energy exchange between vegetation and atmosphere. In his model, Lhomme assumed the sources and sinks of sensible and latent heat fluxes to be uniformally distributed throughout the height of canopy rather than concentrated at a level inside the canopy.

Agricultural crops mostly consist of a large number of roughness elements with irregular shapes being distributed more or less uniformly over some area. The aerodynamic roughness parameter  $(z_0)$ and displacement height (d) are the two important elements of the surface properties that control the surface–atmosphere interactions. Usually,  $z_0$  is interpreted as a length scale that characterises the efficiency for removing momentum from the flow and d is interpreted as an effective level of underlying surface on which source and sink are supposed to lie.

Keywords. Latent heat flux; sensible heat flux; net radiation; aerodynamical resistance; boundary layer resistance; stomatal resistance.

Proc. Indian Acad. Sci. (Earth Planet. Sci.), 110, No. 2, June 2001, pp. 179–184© Printed in India.

The purpose of this paper is to redefine Lhomme's model (1988a) in the frame of roughness length parameter  $(z_0)$  and the zero-displacement height (d). To handle the model more efficiently, the micrometeorological and physiological profiles of the crop have been introduced as initial conditions to arrive at final solutions. The solutions in compact form provide simultaneous calculation algorithms to calculate sensible and latent heat flux densities. The methods of numerical simulation of Lhomme (1988b) have been adopted to operate the model for a standard canopy like a maize crop. The latent heat fluxes computed from reformulated model as functions of stomatal resistance profile and different values of soil surface resistance are presented and compared with those computed by Lhomme (1988b).

#### 2. Basic equations in multi-layer model

The crop canopy, assumed to be horizontally homogeneous, is divided into several parallel thin layers. Subscript *i* refers to the layer number counted from 1 to *n* from top of canopy to soil surface.  $LAI_i$  is the leaf area index of layer *i* per unit ground surface area and  $T_{L,i}$  is the mean temperature of the leaves in the layer *i*.  $e_s(T_{L,i})$  is the vapour pressure inside the stomatal cavity that is assumed to be saturated at leaves temperature  $T_{L,i}$  in layer *i*.  $T_{a,i}$  and  $e_{a,i}$ are the mean temperature and vapour pressure of air in layer *i*.

The Lhomme's model (1988a) is based on electrical analogue in which sensible and latent heat fluxes replace the current, and corresponding potentials are, respectively,  $\rho c_p T$  for sensible heat and  $(\rho c_p / \gamma) e$  for latent heat. In the diffusion process between leaves and air, the latent heat experiences two kinds of resistance: the stomatal resistance  $rs_i$ , and boundary layer resistance  $rb_i$  while the sensible heat flux experiences only boundary layer resistance, assumed to be the same for both transfers. Assuming the atmosphere to be neutral, the elementary fluxes in each layer can be written as:

$$\delta H_i = \rho c_p (T_{L,i} - T_{a,i}) / r e_{c,i}, \qquad (3)$$

$$\delta\lambda E_i = (\rho c_p / \gamma) (e_{s,i}(T_{L,i}) - e_{a,i}) / r e_{\nu,i} \qquad (4)$$

with

$$re_{c,i} = rb_i/2LAI_i,\tag{5}$$

$$re_{\nu,i} = (rs_i + rb_i)/2LAI_i.$$
 (6)

Vertical fluxes denoted by  $H_i$  and  $\lambda E_i$  experience a diffusive resistance  $ra_i$  while crossing the layer *i*. They can be written as:

$$H_i = \frac{\rho c_p (T_{L,i} - T_{a,i})}{r a_{i-1}},$$
(7)

$$\lambda E = \left(\rho c_p / \gamma\right) \left[ \frac{e_{a,i}(T_{L,i}) - e_{a,i}}{r a_{i-1}} \right],\tag{8}$$

where the diffusive resistance  $ra_i$  is related to the eddies diffusivity K(z) within the canopy by the relation:

$$ra_i = \int_{z_i}^{z_{i-1}} \frac{dz}{K(z)}.$$
(9)

# 3. Micrometeorological and physiological profiles

In this section, the micrometeorological and physiological profiles of the canopy are stated and will be used in simultaneity with the basic equations mentioned above.

#### 3.1 Net radiation

The extinction of net radiation within the canopy can be described by Beer's law:

$$Rn = Rn(h) \exp\left[-\alpha_r LAI(z/h)\right], \qquad (10)$$

where h is the height of canopy from the soil surface. The extinction coefficient  $\alpha_r$  depends upon the structure of canopy.  $\alpha_r$  varies from 0.45 to 0.65 for maize and rice crops (Monteith 1976).

The soil flux S is generally taken as a fraction of net radiation reaching the ground and can be expressed as:

$$S = \mu R n \tag{11}$$

with  $\mu = 0.1$  (Campbell 1977).

### 3.2 Aerodynamic resistance

In the neutral condition, assuming the roughness lengths for heat and momentum to be equal, the aerodynamic resistance  $ra_0$  above the top of canopy (reference height) can be expressed as:

$$ra_0 = \frac{\ln\left[(z_r - d)/(h - d)\right]\ln\left[(z_r - d)/z_0\right]}{k(u)^2}, \quad (12)$$

where (u) is the wind speed measured at reference height  $z_r.k$  is the von Karman constant (0.41). The analysis of wind records obtained in near neutral condition showed that means values of d and  $z_0$ for maize and rice crops increase with increasing stand height. For rice and maize crops this dependence can be approximated by (Monteith 1976):

$$d = 1.04h^{0.88}$$
 and  $z_0 = 0.062h^{1.08}$ . (13)

#### 3.3 Boundary layer and stomatal resistances

The boundary layer resistance of the leaves  $rb_i$  in the layer *i* can be related to local wind speed  $u_i$  by the relation (Perrier 1976):

$$rb_i(z) = rb_0 u_i^a(z), \tag{14}$$

where  $rb_n = 50$  (for soil surface) and a = -0.5.

A simple parameterization as a function of global radiation can be used to describe stomatal resistance profile:

$$rs = k_0 / Rg(z), \tag{15}$$

where Rg(z) is the short-wave global radiation given by

$$Rg(z) = Rg(h) \exp\left[-\alpha_r LAI(z/h)\right].$$
(16)

The extinction coefficient  $\alpha_r$  is the same for both; the net radiation and global radiation.  $k_0$  is a parameter which varies as a function of water status. For a completely wet crop,  $k_0$  is zero but equal to  $9 \times 10^5$  for an important water stress.

In view of the above relations, the quantities  $re_{c,i}$ and  $re_{v,i}$  appearing in equations (3), (4), (5) and (6) can be written as:

$$re_{c,i} = rb_i u_i^a(z)/2LAI_i, \qquad (17)$$

$$re_{v,i} = \frac{|k_0 \exp(\alpha_r LAI_i(z_i/h))/Rg_i}{+ rb_i u_i^a(z)|/2LAI_i.}$$
(18)

#### 3.4 Wind speed and eddies diffusivity profiles

Many workers have used K-theory (flux gradient theory) to simulate exchange between vegetation canopy and the atmosphere. For a neutral atmosphere, the wind speed and eddy diffusivity can be assumed to decrease exponentially within the canopy (Choudhary and Monteith 1988; Shuttleworth and Gurney 1990):

$$u(z) = u(h) \exp[-\alpha_w(z/h)], \qquad (19)$$

$$K(z) = K(h) \exp[-\alpha_w(z/h)].$$
(20)

The value of  $\alpha_w$  for maize crops ranges from 2.5 to 3.0 (Monteith 1976). From K-theory, eddy diffusivity K(h) and wind speed u(h) at the canopy level are related as:

$$K(h) = K(0)u(h), \tag{21}$$

where K(0) represents eddy diffusivity at the reference height  $z_r$  given by

$$K(0) = k^{2}(h-d) / \ln[(h-d)/z_{0}].$$
 (22)

The wind speed at canopy level u(h) can be calculated from wind speed (u) measured at the reference height  $z_r$  by the following relation:

$$u(h) = \left[ \ln\{(h-d)/z_0\} / \ln\{(z_r - d)/z_0\} \right](u).$$
(23)

Using equations (19)–(23), integration of (9) yields the following expression for aerodynamic resistance  $ra_i$  within the canopy:

$$ra_{i} = h[1 - \exp(\alpha_{w}\Delta z/h)]$$
$$\exp(\alpha_{w}z_{i-1}/h)/[K(0)\alpha_{w}u(h)], \qquad (24)$$

where  $\Delta z = z_{i-} z_{i-1}$ .

# 4. Lhomme's model

The total fluxes at the top of the canopy can be expressed as the algebraic sum of the contributions of each layer:

$$H_0 = \sum \delta H_i, \ \lambda E_0 = \sum \delta \lambda E_i, \ Rn_0 = \sum \delta Rn_i,$$
(25)

followed by conservation equations:

$$H_{i} = H_{i+1} + \delta H_{i}, \lambda E_{i} = \lambda E_{i+1} + \delta \lambda E_{i}, \text{ and} Rn_{i} = Rn_{i+1} + \delta Rn_{i}.$$
(26)

From equation (10)

$$\delta Rn_i = Rn_i - Rn_{i+1}$$
  
=  $Rn_i [1 - \exp(-\alpha_r LAI_i z_i/h)].$  (27)

The net radiation absorbed in each layer balances convective fluxes of sensible and latent heat:

$$\delta Rn_i = \delta H_i + \delta \lambda E_i. \tag{28}$$

This equation is still valid for layer n (soil surface) if  $\delta Rn_i$  is replaced by  $\delta Rn_i - S$ . Thus the equation (28) can be summed up from 1 to n to give:

$$\sum_{i=1}^{n} \delta R n_i - S = H_0 + \lambda E_0. \tag{29}$$

The equations (11) and (29) yield:

$$\sum_{i=1}^{n} \delta R n_i - S = (1-\mu)$$
$$\times \sum_{i=1}^{n} R n_i [1 - \exp(-\alpha_r L A I_i z_i/h)]. \quad (30)$$

Now, without going into details about the derivations of Lhomme's model (1988a), the final expressions of Lhomme's model for sensible and latent heat fluxes are written as functions of quantities enlisted above.

Linearising the saturated vapour pressure versus temperature curve between  $T_{L,i}$  and  $T_{a,i}$  by the slope  $\Delta$  of the curve determined at air temperature  $T_{a,i}$  at the reference height, we get

$$\Delta = \left[ e_s(T_{L,i}) - e_s(T_{a,i}) \right] / (T_{L,i} - T_{a,i}).$$
(31)

The vapour pressure deficit in each layer is written where as:

$$Da_i = e_s(T_{a,i}) - e_{a,i}.$$
 (32)

Now, the equation (4) can be written as:

$$\delta\lambda E_i = \left(\frac{\rho c_p}{\gamma}\right) \left[ \Delta (T_{L,i} - T_{a,i}) + Da_i \right] / r e_{\nu,i}.$$
(33)

The equations (8), (27), and (33) together give:

$$T_{L,i} - T_{a,i} = d_i \delta R n_i / \rho c_p - d_i D a_i / \gamma r e_{\nu,i}, \quad (34)$$

where

$$d_{i} = \frac{re_{\nu,i}re_{c,i}}{re_{\nu,i} + (\rho c_{p}/\gamma)re_{c,i}}.$$
(35)

Lhomme (1988a) gave the following expressions for  $H_0$  and  $\lambda E_0$  respectively:

$$H_{0} = \sum_{i=1}^{n} d_{i} \delta R n_{i} / r e_{c,i} - (\rho c_{p} / \gamma)$$
$$\sum_{i=1}^{n} d_{i} D a_{i} / r e_{\nu,i} r e_{c,i}, \qquad (36)$$

$$\lambda E_0 = (\Delta/\gamma) \sum_{i=1}^n d_i \delta R n_i / r e_{\nu,i} + (\rho c_p / \gamma)$$
$$\sum_{i=1}^n d_i D a_i / r e_{\nu,i} r e_{c,i}$$
(37)

with  $Da_i$  in recurrent form

$$Da_{i} = \alpha_{i} Da_{1} + \beta_{i} \Delta J_{0} ra_{1} / \rho c_{p} + \sum_{i=1}^{n} \varepsilon_{i}^{j} \delta Rn_{i} / \rho c_{p}.$$
(38)

The coefficients  $\alpha_i, \beta_i$  and  $\varepsilon_i^j$  can be evaluated from the following relations:

$$\alpha_{i+1} = a_i \alpha_i + b_i \alpha_{i-1},$$
  

$$\beta_{i+1} = a_i \beta_i + b_i \beta_{i-1},$$
  

$$\varepsilon_{i-j}^{j < i-1} = a_i \varepsilon_i^j + b_i \varepsilon_{i-1}^j,$$
  

$$\varepsilon_{i+1}^{i-1} = a_i \varepsilon_i^{i-1} = a_i c_{i-1},$$
  

$$\varepsilon_{i+1}^i = c_i$$
  
(39)

with first coefficients defined as:

$$\alpha_1 = 1, \beta_1 = 0, \alpha_2 = a_1, \beta_2 = 1, \varepsilon_2^1 = c_1,$$
 (40)

$$b_{i} = -\frac{ra_{i}}{ra_{i-1}},$$

$$a_{i} = 1 - b_{i} - c_{\nu,i} + (\Delta/\gamma)(c_{\nu,i} - c_{c,i})d_{i}/re_{\nu,i},$$

$$c_{c,i} = -\frac{ra_{i}}{re_{c,i}},$$

$$c_{\nu,i} = -\frac{ra_{i}}{re_{\nu,i}},$$

$$c_{i} = \Delta(c_{c,i} - c_{\nu,i}),$$

$$J_{0} = H_{0} - (\gamma/\Delta)\lambda E_{0}.$$
(41)

Lhomme (1988a) further simplified the expressions using Monteith equations (1981) for the total flux density at the top of the canopy:

$$\lambda E_0 = \left[ \Delta (Rn - S) + \rho c_p (Da_0 - Da_1) / ra_0 \right] / (\Delta + \gamma), \quad (42)$$

where  $Da_0$  is the saturation deficit of air at reference height above the canopy and  $ra_0$  is the aerodynamic resistance between the reference height and canopy level.  $ra_0$  can be determined from the equation (12).

To keep consistency with equation (42), Lhomme (1988a) expressed  $Da_1$  as a function of  $Da_0$  such as:

$$Da_1 = Da_0 + \alpha_1 r a_0 \Delta J_0 / \rho c_p. \tag{43}$$

Inserting  $Da_1$  in relation (38) and then substituting  $Da_i$  in relations (36), (37) and again making use of (29) and (30), the final expressions for redefined form of Lhomme's model (1988a) are obtained as:

$$H_{0} = \left[\gamma(1+A+B)(1-\mu)\sum_{i=1}^{n} Rn_{i}\{1-\exp(-\alpha_{r}LAI_{i}z_{i}/h)\} - \sum_{i=1}^{n} E_{i}Rn_{i}\{1-\exp(-\alpha_{r}LAI_{i}z_{i}/h)\} - \rho c_{p}ADa_{0}/ra_{0}\right]/$$

$$\{\gamma + (\Delta + \gamma)(A+B)\}, \qquad (44)$$

$$\lambda E_{0} = \left[\gamma(A+B)(1-\mu)\sum_{i=1}^{n} Rn_{i}\{1-\exp(-\alpha_{r}LAI_{i}z_{i}/h)\} + \sum_{i=1}^{n} E_{i}Rn_{i}\{1-\exp(-\alpha_{r}LAI_{i}z_{i}/h)\} + \sum_{i=1}^{n} E_{i}Rn_{i}\{1-\exp(-\alpha_{r}LAI_{i}z_{$$

$$(-\alpha_r LAI_i z_i/h) \} + \rho c_p ADa_0/ra_0 ] / \{\gamma + (\Delta + \gamma)(A + B) \},$$
(45)

Table 1. For a wet soil surface  $(rs_n = 0)$ , the computed  $\lambda E_0(Wm^{-2})$  are shown as functions of stomatal resistance profile specified by the value of  $k_0 \cdot \lambda E_0^*$  is denoted for Lhomme's model.

$k_0(10^5)$	0	1	2	3	4	5	6	7	8	9
$\lambda E_0 \ \lambda E_0^*$	$586 \\ 564$	$\begin{array}{c} 422 \\ 409 \end{array}$	$339 \\ 330$	289 281	$256 \\ 249$	$231 \\ 226$	$\begin{array}{c} 211 \\ 208 \end{array}$	$\begin{array}{c} 195 \\ 194 \end{array}$	181 183	$\begin{array}{c} 171 \\ 174 \end{array}$

Table 2. For dry soil surface, the same variations as in table 1. The soil surface resistance is taken equal to the stomatal resistance of the last vegetation layer.

$k_0(10^5)$	0	1	2	3	4	5	6	7	8	9
$\lambda E_0 \ \lambda E_0^*$	$586 \\ 564$	$402 \\ 391$	$\begin{array}{c} 307\\ 300 \end{array}$	$\begin{array}{c} 248 \\ 243 \end{array}$	$\begin{array}{c} 206 \\ 204 \end{array}$	$\begin{array}{c} 179\\176 \end{array}$	$\begin{array}{c} 154 \\ 155 \end{array}$	$\begin{array}{c} 139\\ 138 \end{array}$	$\begin{array}{c} 123 \\ 125 \end{array}$	111 114

where

$$\varepsilon_{i}^{i} = \Delta r e_{c,i},$$

$$E_{i} = \sum_{j=1}^{n} d_{j} \varepsilon_{j}^{i} / r e_{c,j} r e_{\nu,j},$$

$$A = r a_{0} \sum_{i=1}^{n} d_{i} \alpha_{i} / r e_{c,i} r e_{\nu,i},$$

$$B = r a_{1} \sum_{i=1}^{n} d_{i} \beta_{i} / r e_{c,i} r e_{\nu,i}.$$
(46)

#### 5. Numerical simulation

All elementary resistances such as stomatal, and boundary layer, are supposed to be known as also the net radiation at canopy level. From the measured height of the canopy, zero-displacement height (d) and roughness length  $(z_0)$  can be calculated (equation 13). Thereafter, the wind speed (u) measured at reference height level  $(z_r)$  can be used to determine the wind speed u(h) at the canopy level (h) from equation (23). The eddies diffusivity K(0) at reference level  $z_r$  calculated from equation (22) will yield eddies diffusivity K(h)at canopy level from equation (21). Similarly, the aerodynamic resistance  $ra_0$  at the reference level can be determined from equation (12). The quantities u(h), K(h) and K(0) already calculated can be inserted in equation (24) to get the aerodynamic resistance  $ra_1$  at the canopy level.

In this way, all quantities computed at the canopy level will be treated as initial values for further computations down to soil level by assigning i = 1. Some computed results are presented for microclimate of a maize crop and compared with those presented by Lhomme (1988b). The physiological characteristics of maize crop (Lhomme 1988b) are as follows:

 $\begin{aligned} \alpha_w &= 2.75, \\ Rn &= 60\% \text{ of Rg.} \end{aligned}$ 

The climatic characteristics at reference height  $(z_r)$  of 3m are:

Air temperature  $(T_{a,0})$  : 25°C, Vapour pressure  $(e_{a,0})$  : 2000Pa, Wind speed (u) :  $3ms^{-1}$ , Global radiation (Rg) :  $800Wm^{-2}$ .

Using the input values mentioned above, computations for latent heat fluxes  $(\lambda E_0)$  have been carried out and the results so found are presented in the tables 1–3. Corresponding latent heat fluxes (denoted by  $\lambda E_0^*$ ) calculated by Lhomme (1988b) are also shown in the tables. In table 1, the soil surface is considered to be wet  $(rs_n = 0)$ . In table 2, the soil surface is considered to be dry so that  $rs_n$  is set equal to the stomatal resistance of the last vegetation layer. In table 3, for a given stomatal profile corresponding to  $k_0 = 4 \times 10^5$ , latent heat fluxes are shown as functions of soil surface resistance  $rs_n$ .

From tables 1 and 2, it is seen that the redefined model yields the value of  $\lambda E_0$  greater than that of  $\lambda E_0^*$  given by Lhomme's model (1988b) but the increase is very much pronounced when the crop canopy is completely wet  $(k_0 = 0)$ . As the stomatal resistance  $(rs_i)$  of the canopy increases,  $\lambda E_0$  gets gradually closer to  $\lambda E_0^*$  and finally becomes almost equal at highest water stress  $(k_0 = 9 \times 10^5)$ . Similar characteristics are seen in table 2 too. In table 3, the trend looks similar. Initially,  $\lambda E_0$  decreases as the soil surface resistance  $(rs_n)$  increases but later

Table 3.  $\lambda E_0$  and  $\lambda E_0^*$  are shown as functions of soil resistance  $(sm^{-1})$  for a given stomatal profile corresponding to  $k_0 = 4 \times 10^5$ . The stomatal resistance of the last vegetation layer is  $2700 sm^{-1}$ .

$rs_n$	0	500	1000	1500	2000	2500	3000	3500	4000	4500	5000	5500	6000	7000
$\begin{array}{c} \lambda E_0 \\ \lambda E_0^* \end{array}$	$256 \\ 249$	$229 \\ 225$	$\begin{array}{c} 214 \\ 216 \end{array}$	$\begin{array}{c} 208 \\ 210 \end{array}$	$\begin{array}{c} 206 \\ 207 \end{array}$	$\begin{array}{c} 204 \\ 205 \end{array}$	$\begin{array}{c} 201 \\ 203 \end{array}$	$\begin{array}{c} 199 \\ 202 \end{array}$	$\begin{array}{c} 198.6\\ 201 \end{array}$	$\begin{array}{c} 198 \\ 200 \end{array}$	197.7	197.2	$\begin{array}{c} 196.4 \\ 199 \end{array}$	$\begin{array}{c} 196 \\ 198 \end{array}$

becomes almost constant as the soil surface resistance goes on increasing.

# 6. Conclusion

Lhomme (1988a) extended mathematically the Penmann's formulae to multi-layer model with expressions partitioning available radiative energy into sensible and latent heat fluxes. Lhomme assumed the sources and sinks to be distributed uniformly throughout the height of the canopy. Further, the model does not seem to be very elaborate in dealing with the micrometeorological and physiological profiles of the crop explicitly.

In this paper, Lhomme's model (1988a) is redefined with the assumption that the sources and sinks lie on a fictitious plane so called zero-displacement plane. In the frame of zerodisplacement height and roughness length, expressions for sensible and latent heat fluxes have been obtained as functions of micrometeorological and physiological profiles of crop. Lhomme (1988b) presented some computed values of latent heat flux for maize crop stand. Using the same input values, latent heat flux has been calculated from redefined model for maize crop. On comparing the results, it is found that the redefined model yields the flux greater than the Lhomme's model (1988b) when the crop is completely wet. On the other hand, the two models are found to be nearly identical at a large value of stomatal resistance of the canopy.

The advantage with the redefined model is that the input parameters are systematically defined, thus giving their initial values at the top of the canopy or at a reference height. There is no need to compute separately the micrometeorological and physiological parameters as in the case of the Lhomme's model. In this way, the redefined model offers simplified calculation algorithms.

# References

- Campbell G S 1977 An introduction to environmental biophysics; (New York: Heidelberg Science Library)
- Chen J 1984 Uncoupled multi-layer model for the transfer of sensible and latent flux densities from vegetation; *Boundary-Layer Meteorology* **26** 213–225
- Choudhary B J and Monteith J L 1988 A four-layer model for budget of homogeneous surfaces; Q. J. R. Meteorol. Soc. 114 373–398
- Deardroff J W 1978 Efficient prediction of ground surface temperature and moisture along with layer of vegetation; J Geophys. Res. 83 1889–1903
- Lhomme J P 1988a Extension of Penmann's formulae to multi-layer model; Boundary-Layer Meteorology 42 281– 291
- Lhomme J P 1988b A generalized combination equation derived from a multi-layer micrometeorological model; Boundary-Layer Meteorology **45** 103–116
- Monteith J L 1976 Vegetation and atmosphere. Vol. 2 Case studies. (London Academic Press) 33–62
- Monteith J L 1981 Evaporation and surface temperature; Q. J. R. Meteorol. Soc. 107–127
- Meyers T and Paw K T 1986 Testing of higher-order closure model for modeling air flow within and above plant canopies; *Boundary-Layer Meteorology* **37** 297–311
- Naot O and Mahrer Y 1989 Modeling microclimateenvironment: A verification study; *Boundary-Layer Mete*orology 46 333–345
- Penmann H L 1948 Natural evaporation from open water, bare soil, and grass; Proc. R. Soc. London A193 120–145.
- Penmann H L 1953 Physical basis of irrigation control; Report 13th International Hort. Cong 2 913–923
- Perrier A 1976 Etude et Essai de Modelisation des Echanges de Masse et d'energie au Niveau des Couverts Vegetaux; *These de Doctorat d'Etat*, (Paris: Universite de Paris 6)
- Shuttleworth W J 1976 A one dimensional theoretical description of vegetation-atmosphere interaction; *Boundary-Layer Meteorology* **10** 273–302
- Shuttleworth W J and Gurney R J 1990 Theoretical relation between foliage temperature and canopy resistance in sparse canopy; Q. J. R. Meteorol. Soc. **116** 497–519
- Waggoner P E and Reifsnyder W E 1968 Simulation of temperature, humidity, and evaporation profiles in a leaf canopy; *J Applied Meteorology* **7** 400–409

MS received 17 September 1998; revised 1 May 2001